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Spectroscopic studies of distorted structure bio-nano molecules

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Abstract

In this paper, we have introduced an algebraic technique to Bio Nano-molecules (Porphyrins & Fullerenes) family to determine the Vibrational spectra. An algebraic model of vibrations of polyatomic BioNano-molecules and present, as an example, the vibrational analysis of stretching modes of Nickel Octaethyl Porphyrin (Ni(OEP)), its Isotopomers and bionano molecules. The algebraic technique obtained the results are compared with experimental data; the results are showing good accuracy. Some reassignments of energy levels and predicts location of states not yet observed.

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1. Introduction

For the last few years, theoretical studies of highly excited vibrational states of polyatomic molecules have been one of the most interesting topics for theoreticians and experimentalists because of the development of new laser spectroscopic techniques. The measurements of highly-excited overtone-combination spectra of molecules have renewed in a theoretical description and understanding of the observed spectral properties. Two approaches have mostly been used so far in an analysis of experimental data: (1) the familiar Dunham like expansion of energy levels in terms of rotations-vibrations quantum numbers and (2) the solution of Schrodinger equation with potentials obtained either by appropriately modifying ab-initio calculations or by more phenomenological methods. In this article, we begin a systematic analysis of overtone-combination spectra and intensities of molecules in terms of novel approach: (3) Vibron model [1-4]. This model is a formulation of the molecular spectral problem in terms of elements of Lie algebra and it contains the same physical information of the Dunham and potential approach.

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However, by making use of the powerful methods of group theory, one is able to obtain the desired results in a much faster and straightforward way. In recent years, these polyatomic bio-molecules (i.e Metalloporphyrins) have numerous importance in the field of Chemical Physics. In case of polyatomic bio- molecules the parameters plays major role in the Vibron model. Of course, we have explicitly described with fewer parameters, the vibrational bands of the triatomic linear molecules HCN, OCS [5] and tetratomic molecules HCCF, HCCD by using an algebraic approach [6]. We have also reported the vibrational bands of tetrahedral molecules CCl_4 , SnBr_4 [7] and polyatomic bio-molecules Nickel Octaethyl porphyrin, Nickel porphyrin molecules using U (2) Vibron model respectively [8]. The advantage of the algebraic approach, as compared to that of Dunham or phonological potential models, is that typically it requires fewer parameters to obtain the same level of accuracy. It also provides a simultaneous description of bending and stretching modes.

We consider in this paper the vibrational modes of polyatomic Bio Molecules i.e Metalloporphyrins and its substituted forms. In Sec. 2, we briefly review the theory of Vibron model of vibrations of polyatomic molecules. In Sec. 3, we describe the stretching vibrations of different Metalloporphyrins and its substitute forms. Conclusions are presented in Sec. 4.

2. Theory of Lie algebraic Methods

A general algebraic model of rotation- vibration spectra of molecules was introduced in and applied to the study of triatomic and tetra-atomic [9] molecules. In this model rotations and vibrations are treated simultaneously. For this reason, it becomes impractical when the number of atoms exceeds four. However, if rotations and vibrations are separated, it is possible to construct a simpler version of the Vibron model, which can be used for vibrational analysis of large molecules. This model, which we call the one-dimensional Vibron model, starts from the observation that the eigen states of a one-dimensional Schrodinger equation with a Morse potential,

$$h(p, x) = p^2/2\mu + D [1 - \exp(-\alpha x)]^2, \quad (1)$$

In the algebraic approach, interaction of the type can be taken into account by introducing two terms, called Casimir and Majorana interactions, respectively. If these are added, the total Hamiltonian becomes,

$$H = E_0 + \sum_{i=1}^n A_i C_i + \sum_{i < j}^n A_{ij} C_{ij} + \sum_{i < j}^n \lambda_{ij} M_{-ij} \quad (2)$$

The operators C_{ij} are diagonal, with matrix elements given by

$$\langle N_i, v_i; N_j, v_j | C_{ij} | N_i, v_i; N_j, v_j \rangle = 4[(v_i + v_j)^2 - (v_i + v_j)(N_i + N_j)] \quad (3)$$

where we have subtracted (relative) the last two terms, since these can be reabsorbed in C_i and C_j . The operators M_{ij} have both diagonal and non-diagonal matrix elements

$$\begin{aligned} \langle N_i, v_i; N_j, v_j | M_{ij} | N_i, v_i; N_j, v_j \rangle &= (N_i v_j + N_j v_i - 2v_i v_j) \\ \langle N_i, v_i + 1; N_j, v_j - 1 | M_{ij} | N_i, v_i; N_j, v_j \rangle &= -[v_j(v_i + 1)(N_i - v_i)(N_j - v_j + 1)]^{1/2} \end{aligned}$$

$$\langle N_i, v_i - 1; N_j, v_j + 1 | M_{ij} | N_i, v_i; N_j, v_j \rangle = - [v_i(v_i + 1)(N_j - v_j)(N_i - v_i + 1)]^{1/2} \quad (4)$$

The eigen values of Eq (4) can be easily evaluated and provide a description of n coupled anharmonic vibrators. The couplings in Eq (4) are only first order, in the sense that the operators M_{ij} annihilate one quantum of vibration in bond j and create one in bond i (or vice-verse). In the same way in which one can improve the description of each bond by considering higher powers in the C-operators, one can also improve the description of the couplings by considering higher order terms in the M-operators. We have derived explicit formulas for couplings in which two quanta are annihilated and one is created, for example,

$$\langle N_i, v_i \pm 2; N_j, v_j \mp 1 | M_{ij}^2 | N_i, v_i; N_j, v_j \rangle \quad (5)$$

and one can generate, in a relatively simple fashion, others, if needed. For applications to the problem described here, these higher order couplings are not needed and we therefore delete them henceforth.

3. Results and discussions

3.1 Sample Vibrational Analysis: Stretching Vibrations of Metalloporphrins and its Isotopomers

As an example of use of the algebraic method we analyze the stretching vibrations of different Metalloporphrins. We number the bonds as shown in figure 1.

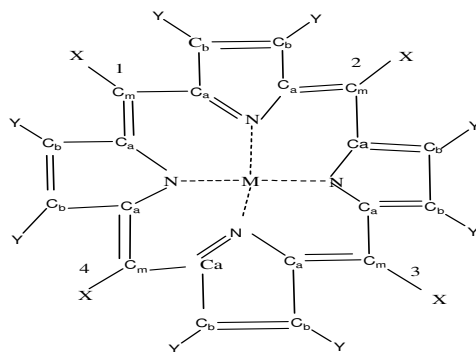


Figure 1. The structure of Metalloporphrins.

Each bond, i , is characterized by its Vibron number N_i , and parameter A_i . The Casimir part of the interbond interactions is characterized by parameter A_{ij} . For the Majorana part we can have, the view of symmetry of the molecule, two possible types of couplings.

One can thus describes all stretching vibrations of **Ni (OEP)-¹⁵N**, **Ni (OEP)-meso-d₄**, **Ni (OEP)-methylene-d₁₆** in terms of five parameters ($N, A, A', \lambda, \lambda'$) as shown in **Table.1**. The Vibron numbers N can be taken from the known anharmonicities of the free C-H or C-C bonds using Eq (3). The data used to determine the parameters are shown in **Table 2**.

Table 1

Comparison between the Observed and Calculated frequencies of the fundamental stretching vibrations

Ni (OEP)- ¹⁵ N				Ni (OEP)-meso-d ₄			Ni (OEP)-methylene-d ₁₆		
Sym	E ^(a) _{obs}	E _{Calc}	Δ	E ^(a) _{obs}	E _{Calc}	Δ	E ^(a) _{obs}	E _{Calc}	Δ
<u>C_m-H vibrational mode</u>									
A _{1g} (ν ₁)	3041	3041.93	-0.93	2263	2258.49	4.51	3044	3042.03	1.97
B _{2g} (ν ₂₇)	3041	3040.89	+0.11	2261	2265.38	-4.38	3041	3040.20	0.80
E _{1u} (ν ₃₆)	3041	3040.01	0.99	2262	2262.30	-0.30	3041	3040.92	0.18
<u>C_b-C_b vibrational mode</u>									
B _{1g} (ν ₂)	1604	1604.04	-0.04	1598	1597.05	0.05	1597	1597.35	-0.35
A _{1g} (ν ₁₁)	1578	1577.96	-0.04	1578	1580.50	-2.50	1567	1568.89	-1.89
E _{1u} (ν ₃₈)	1588	1590.00	-2.00	1586	1587.97	-1.97	1579	1579.05	-0.05

^(a) Observed values taken from the reference [10]**Table 2**Values ^b of Algebraic Parameters Used in the calculation of C_m-H, C_b-C_b Stretching Modes of Porphyrins and its Isotopomers

C _m -H vibrational mode				C _b -C _b vibrational mode		
	Ni (OEP)- ¹⁵ N	Ni (OEP)-meso-d ₄	Ni (OEP)-methylene-d ₁₆	Ni (OEP)- ¹⁵ N	Ni (OEP)-meso-d ₄	Ni (OEP)-methylene-d ₁₆
N	44	44	44	140	140	140
A	-17.6802	-17.6820	-17.614	-2.83	-2.825	-2.835
A'	-0.24	-0.25	-0.28	-1.223	-1.286	-0.452
λ	0.014	0.01136	0.009	0.086	0.068	0.067
λ	0.011	0.012	0.5685	0.047	0.092	0.020

^(b) All values in cm⁻¹ except N, which is dimensionless.

We have used in these tables Wilson numbering for the species of stretching vibrations. It would be very interesting to see whether or not our calculations predict states at the correct. We have explicit calculations up to the fourth overtone (energy up to ≈ 15000 cm⁻¹) as shown in **Table 3 & Table 4**.

Table3.Comparisons between the experimental and Calculated fundamental vibrational frequencies of C₆₀ & C₇₀ (Cm⁻¹).

<u>C₆₀ molecule</u>				
Symmetry	Normal level	Exp ^(a)	Cal	Δ(exp-Cal)
A _{1g}	ν ₁	487	485.9303	-1.0697
	ν ₂	1473	1472.2039	0.7961

T _{1g}	v ₃	561	562.3943	-1.3943
	v ₄	824	826.3922	-2.3922
	v ₅	1277	1277.2034	-0.2034
T _{2g}	v ₆	555	550.2039	4.7961
	v ₇	723	728.3049	-5.3049
	v ₈	789	788.0293	0.9707
G _{1g}	v ₉	1345	1346.3049	-1.3049
	v ₁₀	480	484.3092	-4.3092
	v ₁₁	565	564.8372	0.1628
	v ₁₂	741	741.2345	-0.2345
	v ₁₃	1073	1076.6573	-3.6573
	v ₁₄	1309	1310.7748	-1.7748
H _g	v ₁₅	1507	1510.6253	-3.6253
	v ₁₆	261	266.3948	-5.3948
	v ₁₇	429	430.8272	-1.8272
	v ₁₈	705	704.2812	0.7188
	v ₁₉	772	776.0893	-4.0893
	v ₂₀	1105	1109.2034	-4.2034
	v ₂₁	1251	1260.6592	-9.6592
	v ₂₂	1426	1427.0983	-1.0983
A _u	v ₂₃	1585	1587.2039	-2.2039
	v ₂₄	947	948.3892	-1.3892
T _{1u}	v ₂₅	527	528.2039	-1.2039
	v ₂₆	577	577.9403	-0.9403
	v ₂₇	1190	1191.2389	-1.2389
T _{2u}	v ₂₈	1431	1432.2930	-1.2930
	v ₂₉	336	339.4857	-3.4857
	v ₃₀	709	708.3849	0.6151
	v ₃₁	958	948.3049	9.6951
	v ₃₂	1177	1182.2784	-5.2784
	v ₃₃	1535	1536.2038	-1.2038
G _u	v ₃₄	351	352.3042	-1.3042
	v ₃₅	737	737.3048	-0.3048
	v ₃₆	751	751.2030	-0.2030
	v ₃₇	1308	1310.2933	-2.2933
H _u	v ₃₈	1435	1438.2039	-3.2039
	v ₃₉	400	399.9483	0.0517
	v ₄₀	530	535.4948	-5.4948
	v ₄₁	666	668.3049	-2.3049
	v ₄₂	728	732.2039	-4.2039
	v ₄₃	1220	1220.2039	-0.2039
	v ₄₄	1434	1437.3948	-3.3948
	v ₄₅	1576	1579.2039	-3.2039

Δ(r.m.s) = 78.0888 cm⁻¹

Symmetry	Normal level	<u>C₇₀ molecules</u> Exp ^(a)	Cal	Δ(exp-Cal)
A ₁	v ₁	428	431.9302	-3.9302

A ₂	v ₂	1472	1471.2735	0.7265
	v ₃	629	630.3940	-1.3940
	v ₄	936	940.3048	-4.3048
	v ₅	1215	1210.2934	4.7066
E ₁	v ₆	533	533.4758	-0.4758
	v ₇	722	723.6785	-1.6785
	v ₈	767	767.0920	-0.0920
	v ₉	1350	1352.7839	-2.7839
E ₂	v ₁₀	479	480.4637	-1.4637
	v ₁₁	675	675.0493	-0.0493
	v ₁₂	735	736.4072	-1.4072
	v ₁₃	1055	1050.8373	4.1627
	v ₁₄	1313	1313.7652	0.7652
	v ₁₅	1516	1514.5869	1.4131
	v ₁₆	245	245.7820	-0.7820
	v ₁₇	408	410.2875	-2.2875
	v ₁₈	714	715.8309	-1.8309
	v ₁₉	1169	1167.0392	1.9608
	v ₂₀	1226	1226.5768	-0.5768
	v ₂₁	1257	1260.6592	-3.6592
	v ₂₂	1433	1429.0983	3.9170
	v ₂₃	1574	1575.2039	-1.2039
A ₂ ''	v ₂₄	892	892.3049	-0.3049
A ₁ ''	v ₂₅	507	504.5068	2.4932
	v ₂₆	573	577.3049	-4.3049
	v ₂₇	1177	1181.6099	-4.6099
	v ₂₈	1432	1432.2930	-0.2930
E ₁ '	v ₂₉	438	439.3049	-1.3049
	v ₃₀	564	565.4054	-1.4054
	v ₃₁	1143	1145.5748	-2.5748
	v ₃₂	1463	1465.4950	-2.4950
	v ₃₃	318	316.2039	1.7961
	v ₃₄	704	706.9857	-2.9857
	v ₃₅	896	900.6574	-4.6574
	v ₃₆	1206	1210.0398	-6.0398
	v ₃₇	1568	1570.3049	-2.3049
	v ₃₈	1577	1580.7812	-3.7812
E ₂ '	v ₃₉	304	309.7856	-5.7856
	v ₄₀	701	705.8517	-4.8517
	v ₄₁	920	923.7869	-3.7869
	v ₄₂	1156	1160.6530	-4.6530
	v ₄₃	1517	1520.6652	-3.6652
	v ₄₄	383	384.7854	-1.7854
	v ₄₅	728	728.6755	-0.6755

$$\Delta(\text{r.m.s}) = 68.2436 \text{ cm}^{-1}$$

^(a) Experimental data has taken from the reference[20]

Table 4.

The fitting parameters used in the study of vibrational spectra of fullerene C₆₀ & C₇₀ is given in Table 2.

	<u>Vibron number</u>	<u>Algebraic parameters</u>		
	N	A	λ	λ
C ₆₀	140	-1.2039	0.0739	-0.4932
C ₇₀	140	0.8493	0.0348	0.5903

.....
All values in cm⁻¹ except N, which is dimensionless.

We have used the algebraic Hamiltonian to study the highly excited vibrational levels of the molecule **Ni (OEP)-¹⁵N**, **Ni (OEP)-meso-d₄**, **Ni (OEP)-methylene-d₁₆**. The highly excited vibrational levels, calculated by using the algebraic Hamiltonian Eq. (2), are shown in Figs. 2, 3, 4, 5, 6, & 7.

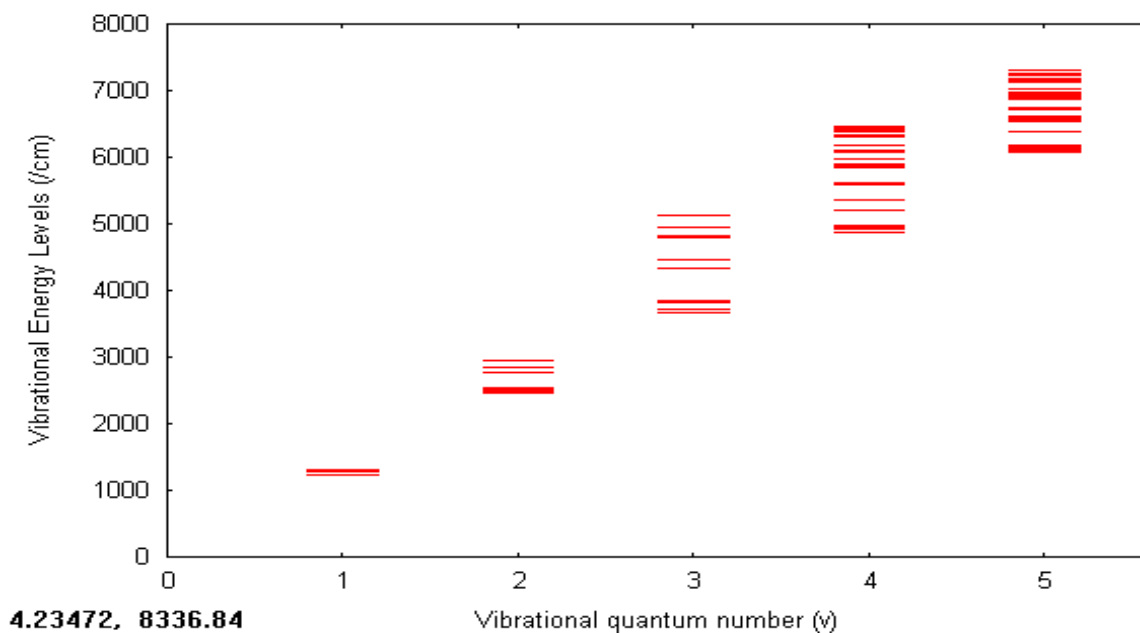
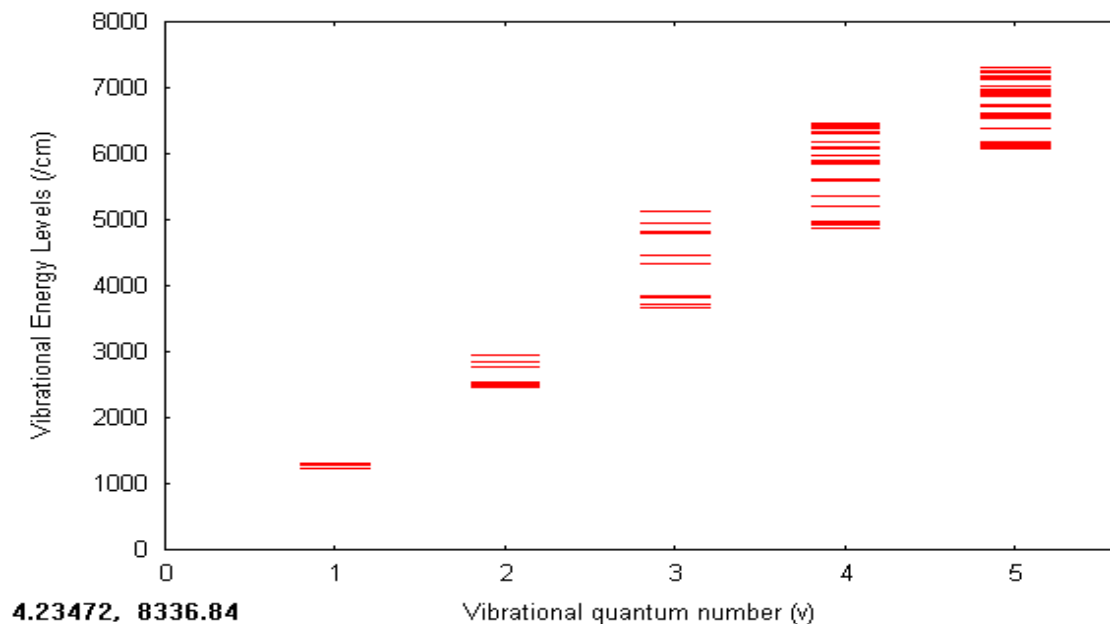
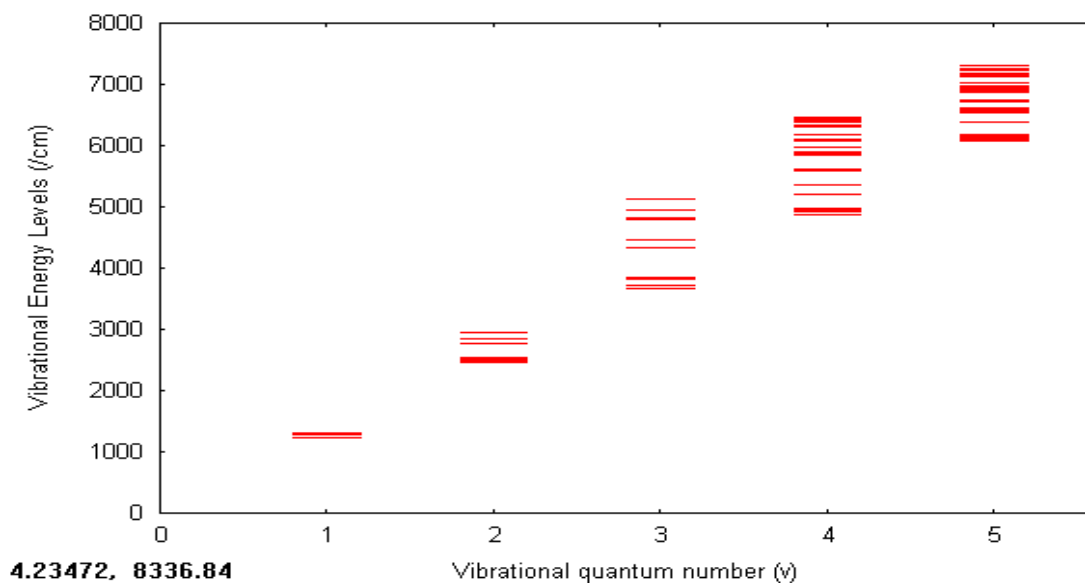
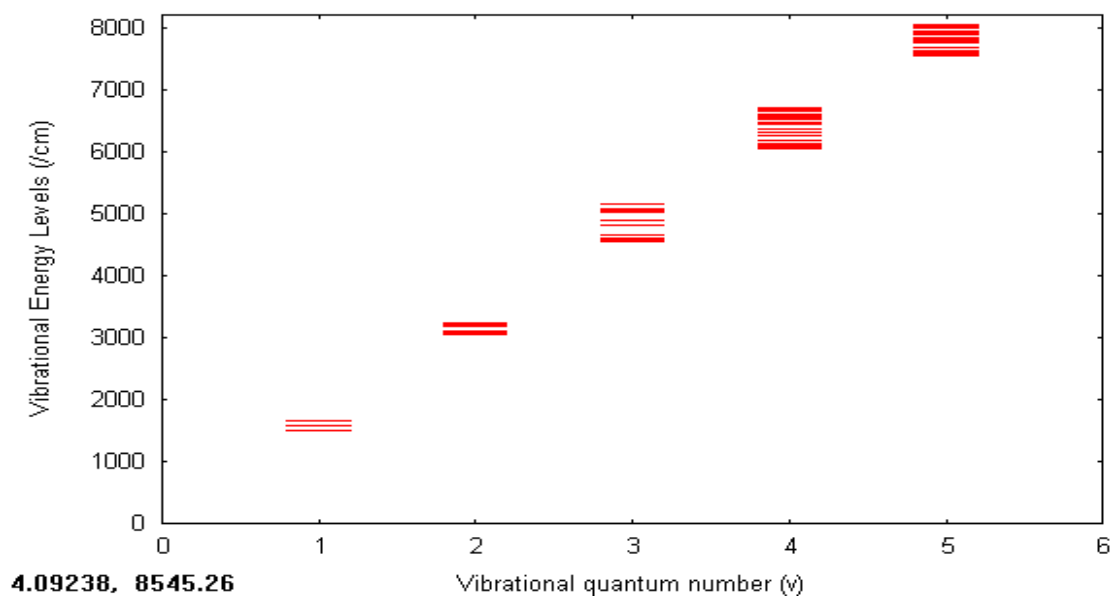
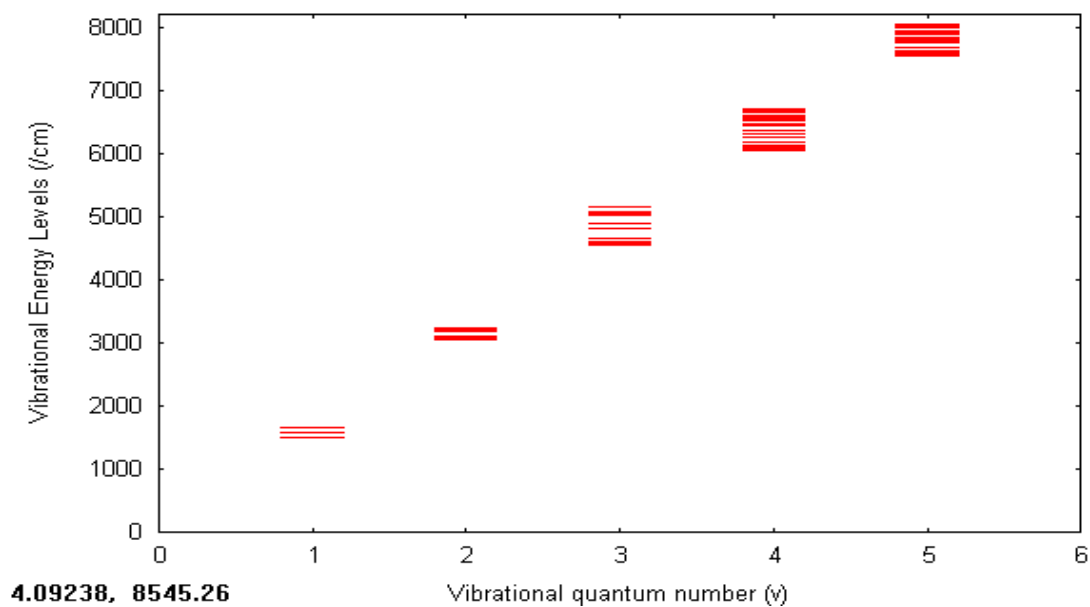


Fig 2. C_m-H band vibrational energy level of **Ni (OEP)-¹⁵N**

Fig 3. C_m-H band vibrational energy level of Ni (OEP)-meso-d₄Fig 4. C_m-H band vibrational energy level of Ni (OEP)-meso-d₄

Fig 5. C_m-H band vibrational energy level of Ni (OEP)-¹⁵NFig 6. C_m-H band vibrational energy level of Ni (OEP)-meso-d₄

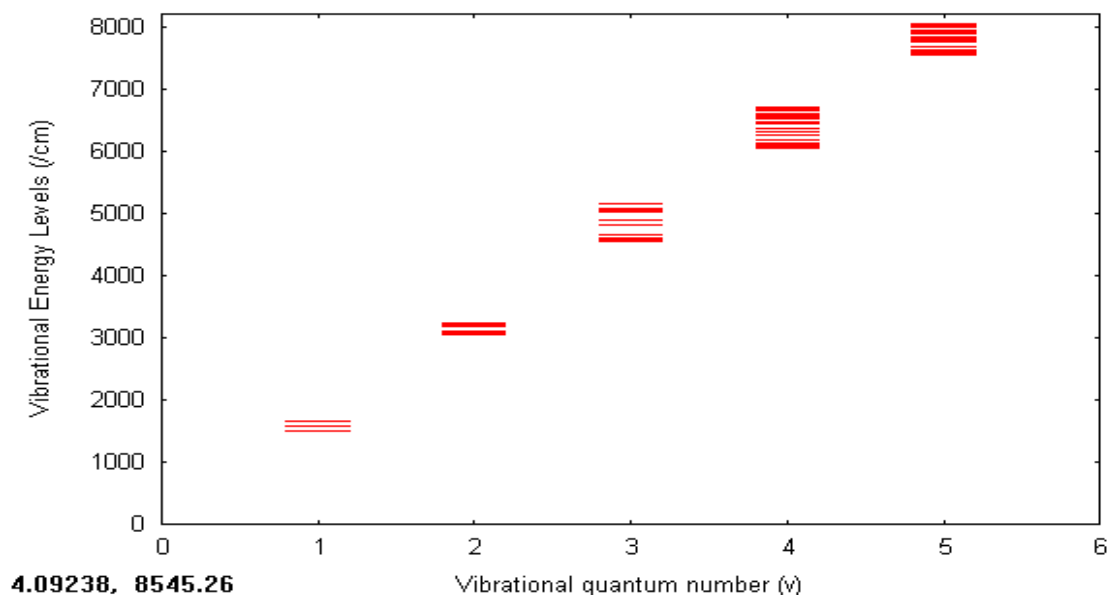


Fig 7. C_m -H band vibrational energy level of Ni (OEP)-meso-d

When the quantum number v increases in a fixed band, the number of energy levels increases rapidly. Usually, the degeneracy or quantidegeneracy of energy levels is called clustering. It may be seen from in Figs. 2, 3,4,5,6 & 7, which the vibrational energy levels of Porphyrins and its Isotopomers make up clusters at excited states.

4. Conclusions

We have presented here a vibrational analysis of the stretching/bending modes of Bio molecules (i.e Nickel Porphyrins) and Nano molecules (Fullerenes C₆₀, C₇₀) in terms of one-dimensional Vibron model i.e U(2) algebraic model. The Nano-molecules C₆₀ and C₇₀ are I_h and D_{5h} point group symmetry respectively.

In this study the vibrational frequencies of Nano molecules C₆₀ and C₇₀ for 7 vibrational bands, we obtain the RMS deviation i.e $\Delta(r.m.s) = 78.0888 \text{ cm}^{-1}$, 68.243 cm^{-1} , and the locality parameters are $\xi_1 = 0.0124$, $\xi_2 = 0.0097$ respectively.

Using improved set of algebraic parameters, the RMS deviation we reported in this study for Bio and Nano molecule is lying near about the experimental accuracy. Using only four algebraic parameters, the RMS deviation we reported in this study for Nano molecule are better fit.

The above two points confirm that in four parameters fit, the set of algebraic parameters we reported in this study of local to normal transition provide the best fit to the spectra of Nano molecules.

We hope that this work will be stimulate further research in analysis of vibrational spectra of other Nano molecules like fullerenes and protein molecules where the algebraic approach has not been applied so far. The research work concern is in progress in case of complex Nano systems.

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